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# Atypical Properties of FIB-Patterned RuO<sub>x</sub> Nanosupercapacitors

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ABSTRACT: Small-scale electrochemical capacitors, also called microsupercapacitors, have been the subject of intense research in the last few years as miniaturized energy storage components for modern electronics. Although numerous microfabrication processes have been successfully assessed to achieve low-profile supercapacitors with sub-millimeter scale features, several advances still need to be made in their performance characteristics to become industrially viable components. Here we report the occurrence of unexpected properties of on-chip supercapacitors when reducing inter-electrode spacing down to the nanometer-scale. An ultrahigh power concomitant with a high capacitance and energy density, an unforeseeable extended cell voltage and an impressive lifetime were obtained at such small dimensions with a FIBpatterned nanosupercapacitor based on RuO<sub>x</sub> pseudocapacitive material. The scaling relationship between miniaturized supercapacitors and electrochemical responses leads to valuable understanding of electrode reactions and rate-limiting steps. This finding offers new opportunities in the design of integrated energy storage devices with improved properties.

While the demand for ever-smaller electronic devices has pushed to the miniaturization of a variety of technologies, energy-storage units have lagged behind in this downsizing tendency.<sup>1</sup> Developing energy-storage elements for self-powered nanosystems, such as nanoelectromechanical systems (NEMS)<sup>2</sup> or wireless nanosensor networks,<sup>3</sup> is of critical importance in many areas of applications, ranging from home automation and industrial tracking to healthcare and environmental monitoring. As the size of power source should be commensurate with the device it powers, great efforts have been recently devoted to integrating energy harvesting<sup>4</sup> and storage<sup>5</sup> nanodevices with the objective to build densely integrated electronic systems having total energy autonomy.

Electrochemical capacitors, also called supercapacitors, are reversible energy-storage devices with a high power density, fast charge and discharge rates, and long service life compared to batteries.<sup>6-8</sup> Nanoscale supercapacitors are consequently of great interest in the development of nanoscale sensor systems requiring the delivery of peak power and long operating lifetime. For nanoscale supercapacitors, it is imperative to consider all reported properties (capacitance, energy, power...) normalized to the footprint area on the chip. However, fabrication of miniaturized planar electrodes at very small scale is a challenge on its own. Optical lithographic efforts have led to the development of microscopic thin-film supercapacitors,<sup>9,10</sup> however the resolution of conventional photolithography techniques is limited by the diffraction of light, which makes it difficult for them to meet the demand of fabricating ultra-high density integrated devices. Actually, most of the "microsupercapacitors" reported in the literature have photolithographic feature size well above 10 µm and typical footprint occupancy of *ca*. 1 cm<sup>2</sup>.<sup>11</sup>

use of focused ion beam (FIB) technology to directly scribe a high-resolution pattern onto a deposited thin film electrode.<sup>12</sup>

Several carbonaceous<sup>13-16</sup> and pseudocapacitive materials<sup>17,18</sup> have been evaluated for miniaturized supercapacitors, considering the choice of materials is one of the factors that ensure high performance. Among them, ruthenium oxide, RuO<sub>x</sub>, appears to be the most promising material owing to its high specific capacitance, highly reversible redox reactions and long cycling life. Even if its high cost has limited its development for large-size supercapacitors, the use of this material becomes relevant for miniaturized supercapacitors where only tiny amount of active material is required.

Here we describe the FIB-fabrication of a nanosupercapacitor based on ruthenium oxide that exhibits singular properties. The manufacturing approach allows the diminution of the cell dimension by six orders of magnitude (from square centimeters to few tens of square micrometers) and the reduction of the inter-electrode spacing by two orders of magnitude (from tens of micrometers to hundreds of nanometers). We show that this scaling down reduces the ohmic losses and improves the energy storage efficiency of the electrochemical devices. We also focus on the anomalous extension of cell voltage observed at very high charge/discharge rate.

A Cr/Pt (40/200 nm) thin film was deposited by evaporation on an oxidized silicon substrate Si/SiO<sub>2</sub> (500 nm), followed by the deposition of a 100 nm-thick metallic Ru by magnetron sputtering under a total pressure of 1 Pa without additional heating.<sup>19</sup> The electrical resistivity of Ru was measured to be *ca*. 40  $\mu$ Ω.cm. Using a suitable lithographic process and an optimized ion beam (Supporting Figure S1 and S2), a narrow gap of 50 nm was created by FIB sputter etching in the thin film of Pt/Ru material with a well-defined pattern, as illustrated in Figure 1. The Ru was then electrochemically oxidized into hydrous ruthenium oxide RuO<sub>x</sub> (Supporting Figure S3)

before being heat-treated at 150°C for 1 h. A protective insulating SU-8 masking layer was subsequently deposited to define an electrochemically active surface exposed to the electrolyte, delimiting a final device area of 25  $\mu$ m<sup>2</sup> (*i.e.* potentially enabling 4 million units per cm<sup>2</sup>).

Figure 2a shows the Nyquist plot of the FIB nanosupercapacitor, with an exceptionally low equivalent series resistance (ESR, Figure 2a inset) of 821  $\Omega_{\mu}$  (*i.e.* 8.2  $\mu\Omega_{\mu}$ cm<sup>2</sup>) and a perfect vertical increase of the impedance in the low-frequency region. The corresponding specific power density is  $304 \,\mu W/\mu m^2$  (*i.e.*  $30.4 \,kW/cm^2$ ) for a cell voltage of 1 V, which is the highest value ever reported for miniaturized on-chip supercapacitors and more than four orders of magnitude higher than state-of-the-art microsupercapacitors.<sup>11,20-23</sup> This prodigious power ability of nanosupercapacitors is explained by the ultra-small physical separation isolating the positive and the negative electrode, which leads to a dramatic reduction of the cell constant and the ESR, as demonstrated in our previous works.<sup>24</sup> However, ultra-high power values should always be treated with caution. There is a trade-off between power and capacitance, and the ultimate performance improvement is to increase the areal power density without sacrificing the areal energy of the device.<sup>9</sup> In the low frequency region, an increase of the imaginary portion of the impedance is observed with a slope close to 90°, as it is expected for capacitive behavior with no leakage current. At such high impedance, the contribution of the dielectric capacitance of the underlying SiO<sub>2</sub> layer,  $C_d$ , cannot be neglected. The measured capacitance, C, is therefore the sum of two capacitances in parallel, the dielectric capacitance and the pseudocapacitance of the nanosupercapacitor,  $C_{sc}$ , so that  $C = C_d + C_{sc}$ . Taking into account  $C_d$ , the cell capacitance deduced from the impedance spectrum at 100 mHz is 180 pF/µm<sup>2</sup> (*i.e.* 18 mF/cm<sup>2</sup>), which is an excellent value compared to those reported for thin-film microsupercapacitors.<sup>9,25</sup> The high specific power of the nanosupercapacitor is thus concomitant with a high specific capacitance

with a very low time constant  $\tau = ESR.C$  of 148 ns. The aptitude to combine high areal energy and high areal power density is also illustrated by the ability of the nanosupercapacitor to maintain a capacitive behavior at both low (Figure 2b) and ultra-high scan rates (Figure 2c). At low scan rates, the cyclic voltammograms (CVs) display the quasi-perfect rectangular shape of an ideal capacitor with no *ESR*. The cell capacitance, estimated from the CV, achieves a reproducible value of 40 pF/ $\mu$ m<sup>2</sup> (*i.e.* 4 mF/cm<sup>2</sup>) when cycled at 5 mV/s, and an excellent areal energy density of 20 pJ/ $\mu$ m<sup>2</sup> (*i.e.* 2 mJ/cm<sup>2</sup>) for a 1 V cell voltage. The volumetric capacitance is estimated to be around 222 pF/ $\mu$ m<sup>3</sup> (*i.e.* 222 F/cm<sup>3</sup>) which is slightly lower to the theoretic value of hydrated ruthenium oxide.<sup>26</sup> The capacitive behavior was also maintained at ultra-high voltage scan rates exceeding 20 V/s with good reversibility, as reported for high-power microsupercapacitors designed to meet the specifications of ac line-filtering applications.<sup>27-30</sup>

Another unexpected consequence of using nanopatterned supercapacitor is the extent of the cell voltage. We evaluated the electrochemical behavior of the RuO<sub>x</sub> based nanosupercapacitor at 10 V/s with increasing voltage window (Figure 3a). Surprisingly, the nanodevice was able to cycle up to 1.7 V in aqueous sulfuric acid electrolyte with no visible sign of degradation due to water electrolysis. This result challenges the long-held axiom that supercapacitors based on aqueous electrolyte have a restricted operating voltage range of *ca*. 1 V because of the electrochemical decomposition of water at higher potentials.<sup>31,32</sup> Moreover, an intriguing phenomenon is that the abrupt current increase related to water splitting appears to move towards higher voltage values with increasing scan rates, from *ca*. 1 V at 100 mV/s up to *ca*. 1.9 V at 50 V/s (Figure 3b). Since the stored energy *E* increases with the square of the voltage ( $E = \frac{1}{2}C.V_0^2$ ), this extension of the nanosupercapacitor voltage leads to a 261 % increase in energy. As a result, the nanosupercapacitor has far more energy when high power delivery is

required in very short times. This behavior is in total disagreement with conventional supercapacitors and microsupercapacitors where the specific energy decreases with increase in specific power because of kinetic and resistance issues.<sup>33</sup>

To understand and analyze this unusual behavior, we have considered the sweep rate dependence of the recorded current at different cell voltages (from 1 to 1.9 V) as shown in Figure 4. A linear dependency between log(current i) and log(scan rate v) is observed for a cell voltage of 1 V in aqueous sulfuric acid electrolyte. The slope of the log(i) vs. log(v) plot is 0.96 on the whole sweep-rate range (from 50 mV/s to 20 V/s), indicating that the overall response tends to be capacitive as expected for a pseudocapacitive material tested within the domain of thermodynamic stability of water. When the nanosupercapacitor is tested at much higher voltages, *i.e.* well above the upper limit for water decomposition (H<sub>2</sub>O) into oxygen (O<sub>2</sub>) and hydrogen (H<sub>2</sub>) gas, an additional faradaic current contribution appears. In Figure 4, this is evidenced by a region at low scan rates where log(i) is independent of scan rate. The faradaic current *i<sub>f</sub>* increases steadily with the cell voltage (from 1.5 to 1.9 V), suggesting the reaction is under kinetic control.<sup>34</sup> Accordingly, for a fixed cell voltage  $V_0 \ge 1.5$  V, the current response is the sum of two contributions, a pseudocapacitive current  $(i_c)$  originating from ruthenium oxide and a faradaic current (*i<sub>f</sub>*) resulting from water electrolysis. For  $V_0 \ge 1.5$  V, we can therefore observe two distinct behaviors according to the sweep-rate: at low and medium scan rates, the faradaic current of water electrolysis prevails on the capacitive current. However, the capacitive current  $i_c = f(v)$  gets more and more important with increasing scan rates v and becomes larger than the faradaic contribution at high scan rates (*i.e.* when  $v > i_f / C$ ). Having a sub microsecond time constant  $\tau$  of 148 ns, the nanosupercapacitors are moreover able to maintain capacitive

behaviors at ultra-high voltage scan rates v.<sup>9</sup> This phenomenon explains the atypical extension of the nanosupercapacitor voltage above 1.5 V.

Furthermore, the ability to operate at such high voltage is illustrated by the remarkable lifetime of the device in harsh conditions – using galvanostatic charge/discharge cycles at 8 nA/ $\mu$ m<sup>2</sup> (*i.e.* 0.8 A/cm<sup>2</sup>) up to 1.5 V. Figure 5 reveals the incredible stability of the device, with no capacitance loss even after ten thousands cycles of charge/discharge. Regarding the galvanostatic cycles, they all exhibit a symmetrical triangular curve with a nearly linear variation of voltage as a function of time at very high current densities, from 2 to 64 nA/ $\mu$ m<sup>2</sup> (*i.e.* 0.2 to 6.4 A/cm<sup>2</sup>) with a high specific energy reaching 130 pJ/ $\mu$ m<sup>2</sup> (*i.e.* 13 mJ/cm<sup>2</sup>) for the cell. The potential IR-drop at the current switching point is also negligible, confirming the extremely small *ESR* of the cell, as aforementioned.

In summary, we have demonstrated that supercapacitors undergo remarkable enhancement in their electrochemical performances when passing from micro to nanodimensions, with unusual and unique properties in terms of specific energy, power, lifetime and cell voltage. Our FIB-patterned nanosupercapacitor displays a cell capacitance of 180 pF/ $\mu$ m<sup>2</sup> (*i.e.* 18 mF/cm<sup>2</sup>) and can deliver an outstanding specific energy per surface area of 325 pJ/ $\mu$ m<sup>2</sup> (32.5 mJ/cm<sup>2</sup>) thanks to a full utilization of the active material involved in the charge storage. The nanosupercapacitor exhibits also an ultralow *ESR* leading to a maximum power density of 304  $\mu$ W/ $\mu$ m<sup>2</sup> (*i.e.* 30.4 kW/cm<sup>2</sup>), which is orders of magnitude higher than any result reported with alternative devices. Electrochemical investigations carried out in 0.5 M H<sub>2</sub>SO<sub>4</sub> solution also confirmed the possibility to extend, under certain conditions, the operating voltage of up to 1.9 V in aqueous solution without any capacitance fade. Although these phenomena arising from nanoengineered supercapacitors have been validated with a RuO<sub>x</sub> pseudocapacitive material in aqueous sulfuric

acid, the extension of the upper voltage limit could also be obtained with carbonaceous electrode materials and organic electrolytes working well above 2.3 V. Consequently, these nanosupercapacitors open up new horizons in the expanding field of integrated nanoelectronics and nanosystems, such as nomad electronics, autonomous wireless sensor networks, biomedical implants, nanorobotics, NEMS, active radiofrequency identification (RF-ID) tags and embedded nanosensors. Moreover, the outstanding peculiarities obtained at very low dimensions can be exploited to design larger energy storage devices with probable improved properties by assembling arrayed nanosupercapacitors on large surfaces. Further improvements in the nanofabrication process with devices patterned in interdigitated configuration should also easily results in nano and microsupercapacitors with even higher areal energy and power.

## FIGURES



**Figure 1.** Design of planar nanosupercapacitors with  $Pt/RuO_x$  electrodes. (a) Schematic representation of the final nanodevice. (b, c) Scanning electron microscope image at different magnifications before the oxidation of Ru and the deposition of the SU-8 photoresist.



**Figure 2.** Electrochemical characterization of the nanosupercapacitor soaked in one drop of de-aerated H<sub>2</sub>SO<sub>4</sub>. (a) The Nyquist plot (inset: high frequency region) displays an exceptionally low *ESR* of 821  $\Omega$ .µm<sup>2</sup> (8.2 µ $\Omega$ .cm<sup>2</sup>) with a vertical increase of the impedance. (b, c) Both CVs obtained at low (b) and ultra-high (c) scan rates exhibit a pronounced capacitive behavior, with a cell capacitance value of 40 pF/µm<sup>2</sup> (4 mF/cm<sup>2</sup>).



**Figure 3.** Anomalous extension of the nanosupercapacitor voltage in aqueous H<sub>2</sub>SO<sub>4</sub> electrolyte. (a) CVs obtained at 10 V/s in 0.5 M H<sub>2</sub>SO<sub>4</sub>. The nanodevice is able to cycle up to 1.7 V with no oxidation/reduction peaks related to water electrolysis. (b) Evolution of the upper limit voltage *vs.* scan rate. The potential threshold moves toward higher values with increasing scan rates.



**Figure 4.** Sweep-rate dependence of the current of the nanosupercapacitor in de-aerated 0.5 M H<sub>2</sub>SO<sub>4</sub> at different cell voltages. At 1 V, the slope of the plot log(i) vs. log(v) is 0.96. Above 1 V, the experimental points are fitted by adding on the whole sweep-rate range a constant value *k* to the capacitive current *i*<sub>c</sub>.



**Figure 5.** Galvanostatic charge/discharge measurements within a cell voltage of 1.5 V. (a) Even at very high current densities, from 2 to 64 nA/ $\mu$ m<sup>2</sup> (0.2 to 6.4 A/cm<sup>2</sup>), the curves exhibit a symmetrical triangular profile with negligible IR-drop. (b) Capacitance retention under the severe operating conditions of 8 nA/ $\mu$ m<sup>2</sup> (0.8 A/cm<sup>2</sup>) and 1.5 V.

## ASSOCIATED CONTENT

### Supporting Information.

Experimental methods, sweep-rate dependence of voltammetric currents, additional figures, and supporting references.

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Author Contributions

D.P., A.F. and B.R. conceived and designed the experiments for the elaboration of the electrochemical nanocapacitors. J.F.P. was involved in the Ru deposition. S.G., A.F. and D.G. established the electrochemical oxidation of the Ru. A.F. was involved in the photolithographic process and B.R. in the FIB patterning. A.F., A.E. and D.P. performed the electrochemical characterizations. D.P. wrote the paper and all authors discussed the results and commented on the manuscript.

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